## Exercise 2.3.7

Consider the following boundary value problem (if necessary, see Section 2.4.1):

$$
\frac{\partial u}{\partial t}=k \frac{\partial^{2} u}{\partial x^{2}} \quad \text { with } \quad \frac{\partial u}{\partial x}(0, t)=0, \frac{\partial u}{\partial x}(L, t)=0, \quad \text { and } \quad u(x, 0)=f(x)
$$

(a) Give a one-sentence physical interpretation of this problem.
(b) Solve by the method of separation of variables. First show that there are no separated solutions which exponentially grow in time. [Hint: The answer is

$$
\left.u(x, t)=A_{0}+\sum_{n=1}^{\infty} A_{n} e^{-\lambda_{n} k t} \cos \frac{n \pi x}{L} .\right]
$$

What is $\lambda_{n}$ ?
(c) Show that the initial condition, $u(x, 0)=f(x)$, is satisfied if

$$
f(x)=A_{0}+\sum_{n=1}^{\infty} A_{n} \cos \frac{n \pi x}{L} .
$$

(d) Using Exercise 2.3.6, solve for $A_{0}$ and $A_{n}(n \geq 1)$.
(e) What happens to the temperature distribution as $t \rightarrow \infty$ ? Show that it approaches the steady-state temperature distribution (see Section 1.4).

## Solution

Part (a)
The PDE is the governing equation for the temperature in a one-dimensional rod that is homogeneous and has constant cross-sectional area. The boundary conditions indicate that the rod is insulated at the $x=0$ and $x=L$ ends. Initially the temperature distribution in the rod is $u(x, 0)=f(x)$.

## Part (b)

The heat equation and its associated boundary conditions are linear and homogeneous, so the method of separation of variables can be applied. Assume a product solution of the form $u(x, t)=X(x) T(t)$ and substitute it into the PDE

$$
\frac{\partial u}{\partial t}=k \frac{\partial^{2} u}{\partial x^{2}} \quad \rightarrow \quad \frac{\partial}{\partial t}[X(x) T(t)]=k \frac{\partial^{2}}{\partial x^{2}}[X(x) T(t)]
$$

and the boundary conditions.

$$
\begin{array}{lllll}
\frac{\partial u}{\partial x}(0, t)=0 & \rightarrow & X^{\prime}(0) T(t)=0 & \rightarrow & X^{\prime}(0)=0 \\
\frac{\partial u}{\partial x}(L, t)=0 & \rightarrow & X^{\prime}(L) T(t)=0 & \rightarrow & X^{\prime}(L)=0
\end{array}
$$

Now separate variables in the PDE.

$$
X \frac{d T}{d t}=k T \frac{d^{2} X}{d x^{2}}
$$

Divide both sides by $k X(x) T(t)$. Note that the final answer for $u$ will be the same regardless which side $k$ is on. Constants are normally grouped with $t$.

$$
\underbrace{\frac{1}{k T} \frac{d T}{d t}}_{\text {function of } t}=\underbrace{\frac{1}{X} \frac{d^{2} X}{d x^{2}}}_{\text {function of } x}
$$

The only way a function of $t$ can be equal to a function of $x$ is if both are equal to a constant $\lambda$.

$$
\frac{1}{k T} \frac{d T}{d t}=\frac{1}{X} \frac{d^{2} X}{d x^{2}}=\lambda
$$

As a result of applying the method of separation of variables, the PDE has reduced to two ODEs - one in $x$ and one in $t$.

$$
\left.\begin{array}{l}
\frac{1}{k T} \frac{d T}{d t}=\lambda \\
\frac{1}{X} \frac{d^{2} X}{d x^{2}}=\lambda
\end{array}\right\}
$$

Values of $\lambda$ that result in nontrivial solutions for $X$ and $T$ are called the eigenvalues, and the solutions themselves are known as the eigenfunctions. Suppose first that $\lambda$ is positive: $\lambda=\alpha^{2}$. The ODE for $X$ becomes

$$
\frac{d^{2} X}{d x^{2}}=\alpha^{2} X
$$

The general solution is written in terms of hyperbolic sine and hyperbolic cosine.

$$
X(x)=C_{1} \cosh \alpha x+C_{2} \sinh \alpha x
$$

Take a derivative with respect to $x$.

$$
X^{\prime}(x)=\alpha\left(C_{1} \sinh \alpha x+C_{2} \cosh \alpha x\right)
$$

Apply the boundary conditions now to determine $C_{1}$ and $C_{2}$.

$$
\begin{aligned}
& X^{\prime}(0)=\alpha\left(C_{2}\right)=0 \\
& X^{\prime}(L)=\alpha\left(C_{1} \sinh \alpha L+C_{2} \cosh \alpha L\right)=0
\end{aligned}
$$

The first equation implies that $C_{2}=0$, so the second equation reduces to $C_{1} \alpha \sinh \alpha L=0$. Because hyperbolic sine is not oscillatory, $C_{1}$ must be zero for the equation to be satisfied. This results in the trivial solution $X(x)=0$, which means there are no positive eigenvalues. Suppose secondly that $\lambda$ is zero: $\lambda=0$. The ODE for $X$ becomes

$$
\frac{d^{2} X}{d x^{2}}=0
$$

The general solution is obtained by integrating both sides with respect to $x$ twice.

$$
\frac{d X}{d x}=C_{3}
$$

Apply the boundary conditions now.

$$
\begin{aligned}
X^{\prime}(0) & =C_{3}=0 \\
X^{\prime}(L) & =C_{3}=0
\end{aligned}
$$

Consequently,

$$
\frac{d X}{d x}=0
$$

Integrate both sides with respect to $x$ once more.

$$
X(x)=C_{4}
$$

Zero is an eigenvalue because $X(x)$ is not zero. The eigenfunction associated with it is $X_{0}(x)=1$. Solve the ODE for $T$ now with $\lambda=0$.

$$
\frac{d T}{d t}=0 \quad \rightarrow \quad T_{0}(t)=\mathrm{constant}
$$

Suppose thirdly that $\lambda$ is negative: $\lambda=-\beta^{2}$. The ODE for $X$ becomes

$$
\frac{d^{2} X}{d x^{2}}=-\beta^{2} X
$$

The general solution is written in terms of sine and cosine.

$$
X(x)=C_{5} \cos \beta x+C_{6} \sin \beta x
$$

Take a derivative of it with respect to $x$.

$$
X^{\prime}(x)=\beta\left(-C_{5} \sin \beta x+C_{6} \cos \beta x\right)
$$

Apply the boundary conditions now to determine $C_{5}$ and $C_{6}$.

$$
\begin{aligned}
& X^{\prime}(0)=\beta\left(C_{6}\right)=0 \\
& X^{\prime}(L)=\beta\left(-C_{5} \sin \beta L+C_{6} \cos \beta L\right)=0
\end{aligned}
$$

The first equation implies that $C_{6}=0$, so the second equation reduces to $-C_{5} \beta \sin \beta L=0$. To avoid the trivial solution, we insist that $C_{5} \neq 0$. Then

$$
\begin{aligned}
-\beta \sin \beta L & =0 \\
\sin \beta L & =0 \\
\beta L & =n \pi, \quad n=1,2, \ldots \\
\beta_{n} & =\frac{n \pi}{L} .
\end{aligned}
$$

There are negative eigenvalues $\lambda=-n^{2} \pi^{2} / L^{2}$, and the eigenfunctions associated with them are

$$
\begin{aligned}
X(x) & =C_{5} \cos \beta x+C_{6} \sin \beta x \\
& =C_{5} \cos \beta x \quad \rightarrow \quad X_{n}(x)=\cos \frac{n \pi x}{L} .
\end{aligned}
$$

$n$ only takes on the values it does because negative integers result in redundant values for $\lambda$. With this formula for $\lambda$, the ODE for $T$ becomes

$$
\frac{1}{k T} \frac{d T}{d t}=-\frac{n^{2} \pi^{2}}{L^{2}}
$$

Multiply both sides by $k T$.

$$
\frac{d T}{d t}=-\frac{k n^{2} \pi^{2}}{L^{2}} T
$$

The general solution is written in terms of the exponential function.

$$
T(t)=C_{7} \exp \left(-\frac{k n^{2} \pi^{2}}{L^{2}} t\right) \quad \rightarrow \quad T_{n}(t)=\exp \left(-\frac{k n^{2} \pi^{2}}{L^{2}} t\right)
$$

According to the principle of superposition, the general solution to the PDE for $u$ is a linear combination of $X_{n}(x) T_{n}(t)$ over all the eigenvalues.

$$
u(x, t)=A_{0}+\sum_{n=1}^{\infty} A_{n} \exp \left(-\frac{k n^{2} \pi^{2}}{L^{2}} t\right) \cos \frac{n \pi x}{L}
$$

## Part (c)

Apply the initial condition $u(x, 0)=f(x)$ to determine $A_{0}$ and $A_{n}$.

$$
\begin{equation*}
u(x, 0)=A_{0}+\sum_{n=1}^{\infty} A_{n} \cos \frac{n \pi x}{L}=f(x) \tag{1}
\end{equation*}
$$

## Part (d)

To find $A_{0}$, integrate both sides of equation (1) with respect to $x$ from 0 to $x$.

$$
\int_{0}^{L}\left(A_{0}+\sum_{n=1}^{\infty} A_{n} \cos \frac{n \pi x}{L}\right) d x=\int_{0}^{L} f(x) d x
$$

Split up the integral on the left into two and bring the constants in front.

$$
A_{0} \int_{0}^{L} d x+\sum_{n=1}^{\infty} A_{n} \underbrace{\int_{0}^{L} \cos \frac{n \pi x}{L} d x}_{=0}=\int_{0}^{L} f(x) d x
$$

Consequently,

$$
\begin{gathered}
A_{0} L=\int_{0}^{L} f(x) d x \\
A_{0}=\frac{1}{L} \int_{0}^{L} f(x) d x
\end{gathered}
$$

To find $A_{n}$, multiply both sides of equation (1) by $\cos (m \pi x / L)$, where $m$ is an integer,

$$
A_{0} \cos \frac{m \pi x}{L}+\sum_{n=1}^{\infty} A_{n} \cos \frac{n \pi x}{L} \cos \frac{m \pi x}{L}=f(x) \cos \frac{m \pi x}{L}
$$

and then integrate both sides with respect to $x$ from 0 to $L$.

$$
\int_{0}^{L}\left(A_{0} \cos \frac{m \pi x}{L}+\sum_{n=1}^{\infty} A_{n} \cos \frac{n \pi x}{L} \cos \frac{m \pi x}{L}\right) d x=\int_{0}^{L} f(x) \cos \frac{m \pi x}{L} d x
$$

Split up the integral on the left into two and bring the constants in front.

$$
\underbrace{A_{0} \int_{0}^{L} \cos \frac{m \pi x}{L} d x}_{=0}+\sum_{n=1}^{\infty} A_{n} \int_{0}^{L} \cos \frac{n \pi x}{L} \cos \frac{m \pi x}{L} d x=\int_{0}^{L} f(x) \cos \frac{m \pi x}{L} d x
$$

Because the cosine functions are orthogonal, the remaining integral on the left is zero if $n \neq m$. As a result, every term in the infinite series vanishes except for the one where $n=m$.

$$
A_{n} \int_{0}^{L} \cos ^{2} \frac{n \pi x}{L} d x=\int_{0}^{L} f(x) \cos \frac{n \pi x}{L} d x
$$

Consequently,

$$
\begin{aligned}
& A_{n}\left(\frac{L}{2}\right)=\int_{0}^{L} f(x) \cos \frac{n \pi x}{L} d x \\
& A_{n}=\frac{2}{L} \int_{0}^{L} f(x) \cos \frac{n \pi x}{L} d x
\end{aligned}
$$

## Part (e)

The temperature distribution approaches equilibrium as $t \rightarrow \infty$.

$$
\begin{aligned}
\lim _{t \rightarrow \infty} u(x, t) & =\lim _{t \rightarrow \infty}\left[A_{0}+\sum_{n=1}^{\infty} A_{n} \exp \left(-\frac{k n^{2} \pi^{2}}{L^{2}} t\right) \cos \frac{n \pi x}{L}\right] \\
& =A_{0} \\
& =\frac{1}{L} \int_{0}^{L} f(x) d x
\end{aligned}
$$

In particular, the equilibrium temperature distribution is the average of the initial temperature distribution.

